

6'

COMPARISON HEAT CONDUCTION IBVP FOR FINITE AND INFINITE DOMAIN. Convolution theorem.

Finite Rod

$$\begin{cases} U_t = K U_{xx}, & -L < x < L \\ U(-L, t) = U(L, t) \\ U_x(-L, t) = U_x(L, t) \\ U(x, 0) = f(x) \end{cases}$$

Soln:

$$\begin{aligned} U(x, t) &= \sum_{n=0}^{\infty} \left[a_n \cos \frac{n\pi}{L} x + b_n \sin \frac{n\pi}{L} x \right] e^{-K \left(\frac{n\pi}{L} \right)^2 t} \\ &= \sum_{n=-\infty}^{\infty} c_n e^{-i \frac{n\pi}{L} x} e^{-K \left(\frac{n\pi}{L} \right)^2 t} = \sum_{n=-\infty}^{\infty} c_n e^{-i \omega_n x} e^{-K \omega_n^2 t} \end{aligned}$$

$\omega_n = \frac{n\pi}{L}$

Where $c_n = \frac{1}{2L} \int_{-L}^L f(x) e^{i \frac{n\pi}{L} x} dx = \frac{1}{2L} \int_{-L}^L e^{i \omega_n x} f(x) dx$

Infinite rod:

$$\begin{cases} U_t = K U_{xx}, & -\infty < x < \infty \\ U(-\infty, t) = 0, & U(\infty, t) = 0 \\ U(x, 0) = f(x) \end{cases}$$

Soln:

$$U(x, t) = \int_{-\infty}^{\infty} C(\omega) e^{-i \omega x} e^{-K \omega^2 t} d\omega \quad (7.1)$$

Where

$$C(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(\bar{x}) e^{i \omega \bar{x}} d\bar{x} \quad (7.2)$$

Solution of Heat Equation on an infinite domain.

$$\begin{cases} u_t = K u_{xx}, & -\infty < x < \infty, \quad t > 0 & (5.0) \\ u(x, 0) = f(x) & & (5.0.1) \\ + \text{B.C's at } \infty : u(-\infty, t) = 0, \quad u(+\infty, t) = 0 & & (5.0.2) \end{cases}$$

Separation of variables leads to

$$u(x, t) = \phi(x) h(t)$$

$$\textcircled{\text{I}} \quad \frac{dh}{dt} + \lambda K h = 0. \Rightarrow h(t) = A e^{-\lambda K t} \quad (5.1)$$

$$\textcircled{\text{II}} \quad \text{EVP: } \begin{cases} \frac{d^2 \phi}{dx^2} + \lambda \phi = 0 & (5.1) \\ |\phi(-\infty)| < \infty, \quad |\phi(+\infty)| < \infty & (5.2) \end{cases}$$

Why not $\phi(+\infty) = 0, \quad \phi(-\infty) = 0$?

In order to satisfy (5.1) and (5.2) $\lambda \geq 0$

(If $\lambda < 0$, $\phi(x)$ would be a combination of exponentials and won't be bounded as $x \rightarrow \infty$ and $x \rightarrow -\infty$ at the same time.)

Thus,
$$\phi(x) = C_1 \cos(\sqrt{\lambda} x) + C_2 \sin(\sqrt{\lambda} x)$$

Satisfies both (5.1) and (5.2) for all $\lambda \geq 0$.

All $\lambda > 0$ are eigenvalues of (5.1) and (5.2)

This is called "Continuous spectrum."

Extending the Superposition principle from the discrete to the continuous case; we obtain

$$u(x,t) = \int_0^{\infty} [C_1(\lambda) \cos(\sqrt{\lambda}x) + C_2(\lambda) \sin(\sqrt{\lambda}x)] e^{-\lambda kt} d\lambda$$

This can be written in terms of ω , if $\lambda = \omega^2$ as

$$u(x,t) = \int_0^{\infty} [A(\omega) \cos \omega x + B(\omega) \sin \omega x] e^{-k\omega^2 t} d\omega \quad (8.1)$$

where

$$A(\omega) d\omega = C_1(\omega^2) 2\omega d\omega = C_1(\lambda) d\lambda$$

$$B(\omega) d\omega = C_2(\omega^2) 2\omega d\omega = C_2(\lambda) d\lambda.$$

Using the I.C., we arrive to

$$f(x) = u(x,0) = \int_0^{\infty} [A(\omega) \cos \omega x + B(\omega) \sin \omega x] d\omega \quad (8.2)$$

(8.1) and (8.2) can be written as (Exercise 10.2.1) ⁹

$$u(x,t) = \int_{-\infty}^{\infty} C(w) e^{-iwx} e^{-kw^2t} dw \quad (9.1)$$

and $f(x) = u(x,0) = \int_{-\infty}^{\infty} C(w) e^{-iwx} dw$

(Inverse Fourier transform of $f(x)$)

If $f(x)$ is absolutely integrable (plays role of orthogonality),

then $C(w) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(\bar{x}) e^{i w \bar{x}} d\bar{x} \quad (9.2)$

(Fourier transform of $f(x)$)

Conclusion:

The Solution of Heat equation

in the infinite line with I.C. $u(x,0) = f(x)$

and BCs: $u(\pm\infty,t) = 0$ is given by (9.1)-(9.2)

10.4 Fourier Transform and Heat Equation.

We will show now that (9.1) and (9.2)

can be combined into a single integral representation given by

$$u(x,t) = \int_{-\infty}^{\infty} f(\bar{x}) \frac{1}{\sqrt{4\pi kt}} e^{-\frac{(x-\bar{x})^2}{4kt}} d\bar{x} \quad (9.3)$$

(9.1) can be expressed as

$$u(x,t) = \int_{-\infty}^{\infty} \left(C(\omega) e^{-k\omega^2 t} \right) e^{-i\omega x} d\omega \quad (10.1)$$

So $u(x,t)$ can be considered the inverse Fourier transform of the product

$$C(\omega) e^{-k\omega^2 t}$$

where $C(\omega)$ ^{(from (9.2))} is the F.T of $f(x)$ (initial condition)

and $e^{-k\omega^2 t}$ is the F.T of $g(x) = \sqrt{\frac{\pi}{kt}} e^{-x^2/4kt}$.

(from Tables or
derivation in Haberman
Appendix 10.3)

For the purpose of reducing
(10.1) to a single integral, we used the
following theorem called
"Convolution Theorem".

Theorem (Convolution Theorem)

If $F(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) e^{i\omega x} dx$, $f(x) = \int_{-\infty}^{\infty} F(\omega) e^{-i\omega x} d\omega$

Fourier transform of $f(x)$ Inverse F.T of $F(\omega)$.

and $G(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} g(x) e^{i\omega x} dx$, $g(x) = \int_{-\infty}^{\infty} G(\omega) e^{-i\omega x} d\omega$

then, the inverse F.T of the product

$$H(\omega) = F(\omega) G(\omega)$$

is $h(x) = \int_{-\infty}^{\infty} F(\omega) G(\omega) e^{-i\omega x} d\omega$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} f(\bar{x}) g(x - \bar{x}) d\bar{x}$$

Proof.

$$h(x) = \int_{-\infty}^{\infty} F(\omega) \left[\frac{1}{2\pi} \int_{-\infty}^{\infty} g(\bar{x}) e^{i\omega \bar{x}} d\bar{x} \right] e^{-i\omega x} d\omega$$

interchanging
order of integration

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} g(\bar{x}) \left[\int_{-\infty}^{\infty} F(\omega) e^{-i\omega (x - \bar{x})} d\omega \right] d\bar{x}$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} g(\bar{x}) f(x - \bar{x}) d\bar{x} \quad \checkmark$$

If we substitute $F(\omega)$ instead of $G(\omega)$, then

$$h(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(\bar{x}) g(x - \bar{x}) d\bar{x} \quad \checkmark$$

Application of convolution theorem.

In (10.1)

$C(\omega)$ is F.T of $f(x)$ (initial cond.)

$e^{-\kappa\omega^2 t}$ is F.T of $g(x) = \sqrt{\frac{\pi}{\kappa t}} e^{-\frac{x^2}{4\kappa t}}$ (From Table)

Then using convolution theorem

$$u(x, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(\bar{x}) \sqrt{\frac{\pi}{\kappa t}} e^{-\frac{(x-\bar{x})^2}{4\kappa t}} d\bar{x}$$

or

$$u(x, t) = \int_{-\infty}^{\infty} f(\bar{x}) \frac{1}{\sqrt{4\pi\kappa t}} e^{-\frac{(x-\bar{x})^2}{4\kappa t}} d\bar{x}$$